

# Automation of multiphysics numerical simulations using open-source software

---

**Jeremy Bleyer**

coll. Paul Bouteiller, Jørgen Dokken, Alice Gribonval, Jack S. Hale, Thomas Helfer, Andrey Latyshev, Corrado Maurini, Maxime Pierre

Ecole Nationale des Ponts et Chaussées, IP Paris



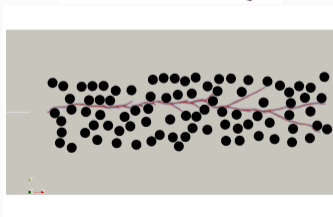
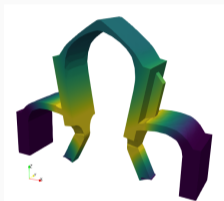
Funded by  
the European Union



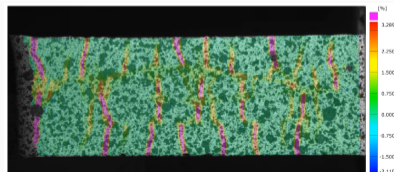
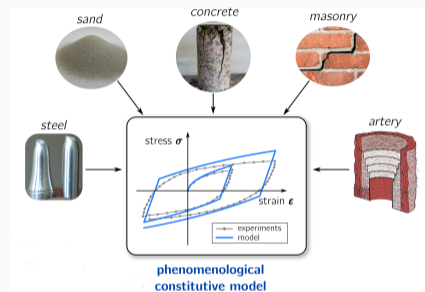
European Research Council  
Established by the European Commission

January 22, 2026

**Structural simulations:** numerical methods (FE), nonlinearities, complex geometry, time-dependency



**Constitutive behavior:** elasticity, viscoelasticity, plasticity, damage, temperature effects...



# Open-source software for computational mechanics

## Advantages/Drawbacks



- code\_aster
- deal\_ii
- elmer
- Kratos

- Cast3m/Manta
- MOOSE
- MFEM

- OpenFOAM
- Basilisk
- FreeFEM++

- gmsh
- pyvista
- Paraview
- pytorch

**Here:** FEniCSx – MFront – JAX

**FEniCSx**

---

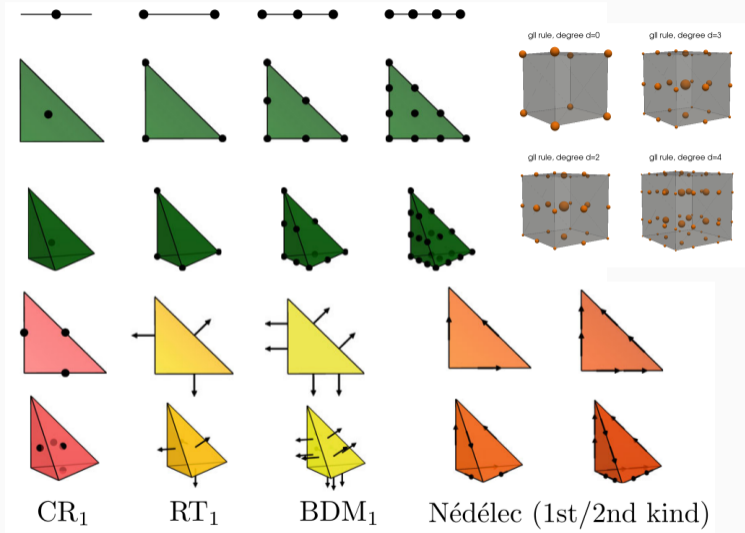
collection of free, open source, software components for **automated solution** of **Partial Differential Equations** (PDEs)



### Features:

- variational formulation expressed with **math symbolic language**: `ufl.grad`, `ufl.div`, `ufl.derivative`, etc.
- **code-generation** (C) and **JIT compilation** of FE kernels
- **extensive library** of finite elements
- designed for **parallel HPC computations**

# Element types



# A simple elasticity example

**Weak form/variational formulation:** Find  $\mathbf{u} \in V$  such that:

$$\int_{\Omega} \boldsymbol{\sigma}(\mathbf{u}) : \nabla^s \mathbf{v} \, d\Omega = \int_{\Omega} \mathbf{f} \cdot \mathbf{v} \, d\Omega \quad \forall \mathbf{v} \in V$$

with **isotropic linear elasticity:**  $\boldsymbol{\sigma}(\mathbf{u}) = \lambda \operatorname{tr}(\boldsymbol{\varepsilon}) \mathbf{I} + 2\mu \boldsymbol{\varepsilon}$  with  $\boldsymbol{\varepsilon} = \nabla^s \mathbf{u}$ .

```
dim, degree = mesh.geometry.dim, 2
V = fem.functionspace(mesh, ("P", degree, dim))

def epsilon(v):
    return ufl.sym(ufl.grad(v))
def sigma(v):
    return lmbda * ufl.tr(epsilon(v)) * ufl.Identity(dim) + 2 * mu * epsilon(v)

u, v = TrialFunction(V), TestFunction(V)

a, L = ufl.inner(sigma(u), epsilon(v)) * dx, ufl.dot(f, v) * dx
```

# Multi-field variational problems

One of the big advantages of **FEniCS**

e.g. **Stokes incompressible fluids**: mixed  $u - p$  formulation

## Taylor-Hood element ( $P_2/P_1$ )

```
# Define function space
P2 = basix.ufl.element('P', tetrahedron, 2, shape=(dim,))
P1 = basix.ufl.element('P', tetrahedron, 1)
The = basix.ufl.mixed_element([P2, P1])
W = fem.functionspace(mesh, The)
```

```
# Define variational problem
(u, p) = TrialFunctions(W)
(v, q) = TestFunctions(W)
a, L = inner(grad(u), grad(v))*dx - p*div(v)*dx + div(u)*q*dx, dot(f, v)*dx
```

# Multi-field variational problems

One of the big advantages of **FEniCS**

e.g. **Stokes incompressible fluids**: mixed  $u - p$  formulation

## MINI element $(P_1 + \mathcal{B})/P_1$

```
# Define function space
Pv1 = basix.ufl.element('P', tetrahedron, 1)
B = basix.ufl.element('Bubble', tetrahedron, 3)
Mini = basix.ufl.mixed_element([Pv1+B, P1])
W = fem.functionspace(mesh, Mini)
```

```
# Define variational problem
(u, p) = TrialFunctions(W)
(v, q) = TestFunctions(W)
a, L = inner(grad(u), grad(v))*dx - p*div(v)*dx + div(u)*q*dx, dot(f, v)*dx
```

## Finite-strain weak equilibrium

$$\int_{\Omega} \mathbf{P}(\mathbf{u}) : \nabla \mathbf{v} \, d\Omega = \int_{\Omega} \mathbf{f} \cdot \mathbf{v} \, d\Omega \quad \forall \mathbf{v} \in V$$

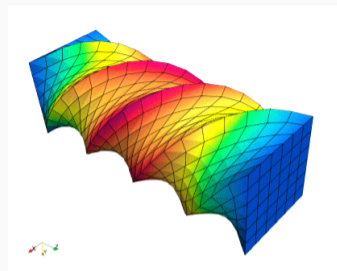
$\mathbf{P} = \frac{\partial \psi}{\partial \mathbf{F}}$ : first Piola-Kirchhoff stress given by a **hyperelastic potential**  $\psi(\mathbf{F})$  with  $\mathbf{F} = \mathbf{I} + \nabla \mathbf{u}$ .

**Example:**  $\psi(\mathbf{F}) = \frac{\mu}{2}(\bar{I}_1 - 3) + k(J^2 + J^{-2} - 2)$  s.t.  $\bar{I}_1 = \text{tr}(\bar{\mathbf{C}})$ ,  $\bar{\mathbf{C}} = J^{-2/3} \mathbf{F}^T \mathbf{F}$ ,  $J = \det \mathbf{F}$

## Tangent operator

$$a_{\text{tangent}}(\mathbf{u}, \mathbf{v}) = \int_{\Omega} \nabla \mathbf{u} : \frac{\partial^2 \psi}{\partial \mathbf{F} \partial \mathbf{F}} : \nabla \mathbf{v} \, d\Omega$$

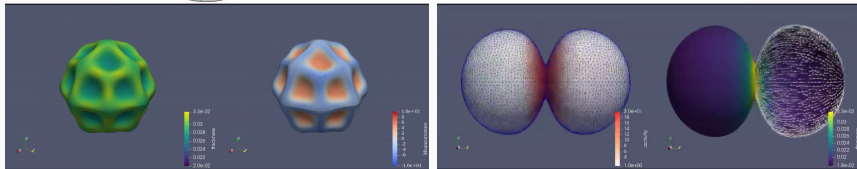
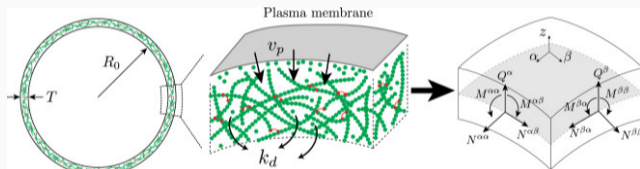
⇒ **UFL symbolic differentiation**



```
def psi(u):  
    ...  
  
E_pot = psi * dx - ufl.dot(f, u)*dx  
Res = ufl.derivative(E_pot, u, v)  
Jac = ufl.derivative(Res, u, du)
```

## Mechanics of biological cells

- The **cell cortex** is a very thin layer under the membrane that controls cell shape.
- It **actively generates forces and flows** (e.g. during cell division or deformation).
- We model its mechanics using an **active viscous shell model** [JMPS, 2022]



**MFront**

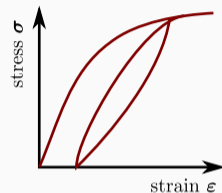
---

# Material constitutive modeling

**Constitutive model** = a **stress-strain** relation

$$\epsilon \longrightarrow \boxed{\text{CONSTITUTIVE RELATION}} \longrightarrow \sigma$$

generally with **custom implementation** in a FE solver

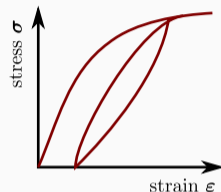


**Constitutive model** = a **stress-strain** relation

$$\epsilon \longrightarrow \boxed{\text{CONSTITUTIVE RELATION}} \longrightarrow \sigma$$

generally with **custom implementation** in a FE solver

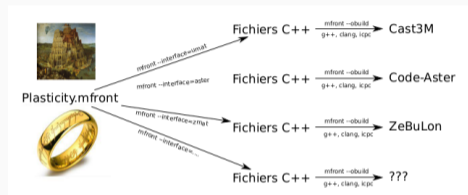
- **decouple** material library from structural solvers
- **inter-operability: material knowledge** can be **shared** across various solvers (industrial and academic)
- improve quality assurance, long-term sustainability, collaboration



MFront = a **code generation tool** dedicated to **material knowledge**

**Main developer:** Thomas Helfer (CEA Cadarache) + many other contributors

- Numerical **efficiency**
- **Inter-operability:** specific (Cast3M, code\_aster, Europlexus, Abaqus, Zebulon, Ansys, etc.) and *generic* interfaces (python, fortran, c, julia, etc.)
- Ease of use (**Domain-Specific Languages**)



## Features

- C++ code, text file input to be compiled for a given interface
- **Typical mechanical behaviors:** elasticity, viscoelasticity, plasticity, damage, temperature dependence, finite strains, cohesive zone models, *homogenized behaviors*
- **solvers:** implicit/explicit time integrators, (quasi)-Newton
- **mtest:** testing tool

## Pre-defined modular bricks

```
@DSL Implicit;
@Behaviour PerfectPlasticity;
@Epsilon 1.e-14;
@Theta 1;
@Brick StandardElastoViscoPlasticity{
  stress_potential : "Hooke" {young_modulus : 200e9,
    poisson_ratio : 0.3},
  inelastic_flow : "Plastic" {
    criterion : "Mises",
    isotropic_hardening : "Linear" {R0 : 150e6},
  }
};
```

## Pre-defined modular bricks

```
@DSL Implicit;
@Behaviour PerfectPlasticity;
@Epsilon 1.e-14;
@Theta 1;
@Brick StandardElastoViscoPlasticity{
  stress_potential : "Hooke" {young_modulus : 200e9,
    poisson_ratio : 0.3},
  inelastic_flow : "Plastic" {
    criterion : "Mises",
    isotropic_hardening : "Linear" {R0 : 150e6},
  }
};
```

```
@DSL Implicit;
@Behaviour Norton;
@Brick StandardElasticity;

@MaterialProperty stress E;
E.setGlossaryName("YoungModulus");
@MaterialProperty real nu, A, nn;
nu.setGlossaryName("PoissonRatio");
A.setEntryName("NortonCoefficient");
nn.setEntryName("NortonExponent");

@StateVariable real p;
p.setGlossaryName("EquivalentViscoplasticStrain");

@Integrator{
  constexpr const auto Me = Stensor4::M();
  const auto mu = computeMu(E, nu);
  const auto sigmae = sigmaeq(sigma);
  const auto ie = 1 / (max(sigmae, real(1.e-12) * E));
  const auto vp = A * pow(sigmae, nn);
  const auto dvp/desigmae = nn * vp * ie;
  const auto n = 3 * deviator(sigma) * (ie / 2);
  // Implicit system
  fe += delta p * n;
  fp -= vp * delta t;
  // jacobian
  dfe/dedelta += 2 * mu * theta * dp * ie * (Me - (n * n));
  dfe/dedelta p = n;
  dfp/dedelta += -2 * mu * theta * dvp/desigmae * delta t * n;
} // end of @Integrator
```

- **Finite strain**: solvers require different "forms" of the tangent operator e.g.  $\frac{d\sigma}{dd}$ ,  $\frac{dP}{dF}$ ,  $\frac{dS}{dE}$  and many others...  $\Rightarrow$  MFfront provides conversion methods
- support for different modeling hypotheses: PlaneStrain, Tridimensional, Axisymmetric, etc. (optimized code/hypothesis)
- **numerical jacobian** (finite-differences) for implicit systems
- compile-time dimensional analysis
- recently, support for generalized behaviors (**multi-physics**)

### Mechanics

$$\Delta \varepsilon_{n+1}, \sigma_n, \alpha_n \rightarrow \boxed{\text{MFfront}} \rightarrow \sigma_{n+1}, \alpha_{n+1}, \frac{d\sigma_{n+1}}{d\varepsilon_{n+1}}$$

### Multiphysics

$$(\Delta \mathbf{g}^1, \dots, \Delta \mathbf{g}^p)_{n+1}, (\sigma^1, \dots, \sigma^p)_n, \alpha_n \rightarrow \boxed{\text{MFfront}} \rightarrow (\sigma^1, \dots, \sigma^p)_{n+1}, \alpha_{n+1}, \frac{\partial \sigma_{n+1}^i}{\partial \mathbf{g}_{n+1}^j}$$

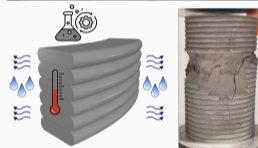
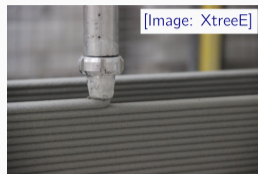
# 3D-printed concrete modeling

Early age concrete = a thermo-poro-**chemo**-mechanical model [M. Pierre et al.]

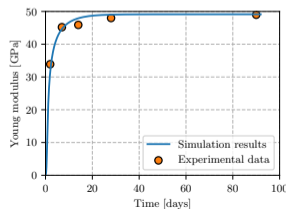
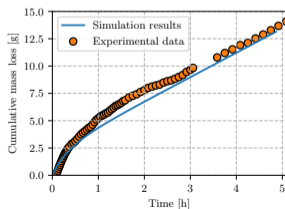
$$\begin{cases} d\boldsymbol{\sigma} = \mathbb{C}(\xi) : d\boldsymbol{\varepsilon} - b(\xi)S_\ell dp\mathbf{1} - 3\alpha K(\xi)dT\mathbf{1} \\ d\phi = b(\xi)\text{tr}(d\boldsymbol{\varepsilon}) + \frac{b(\xi) - \phi_0(\xi)}{K_s} S_\ell dp - 3\alpha (b(\xi) - \phi_0(\xi)) dT - \Delta \bar{V}_s d\xi \\ dS_s = 3\alpha K(\xi)\text{tr}(d\boldsymbol{\varepsilon}) - 3\alpha (b(\xi) - \phi_0(\xi)) S_\ell dp + \frac{C(1 - \phi_0(\xi))}{T_0} dT - \frac{C}{T_0} d\xi \end{cases}$$

Evolution of material properties with hydration, **change of solid volume + heat** due to **chemical reaction(s)**

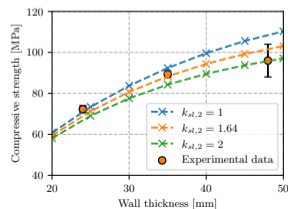
[A. Gribonval et al.]: influence of **environmental conditions** on **compressive strength**



3D-printed concrete

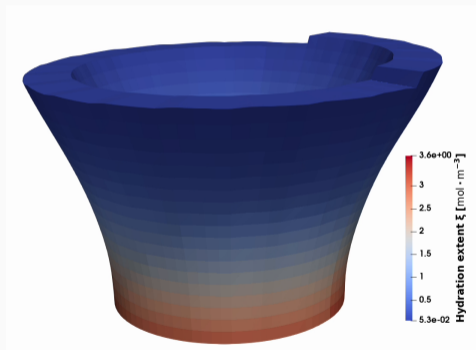


Jeremy Bleyer



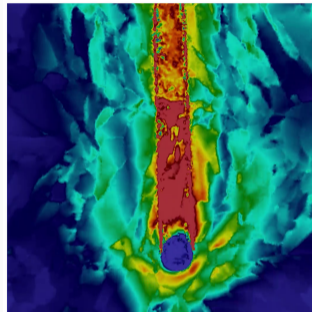
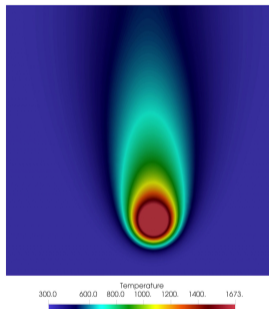
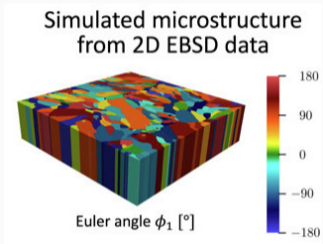
# 3D-printing concrete modeling

[M. Pierre]: gmsh - FEniCSx - MFront pipeline with mesh activation  
THCM model with elastoplastic Cam-Clay behavior



# Residual stresses after laser scanning of additively manufactured stainless steel

coll. Nikhil Mohanan, Manas Upadhyay (LMS, Ecole Polytechnique)



MFront: thermo elastoviscoplastic polycrystal model  
FEniCSx: weakly coupled thermomechanics on polycrystal mesh

**JAX**

---

## Constitutive models as ML models

Constitutive models can be formalized as **parametrized functions** (stress update):

$$\sigma_{n+1}, \alpha_{n+1} = \mathcal{M}_{\theta}(\varepsilon_{n+1}, \alpha_n)$$

where  $\theta$  are the model **material parameters** and  $\mathcal{M}_{\theta} = \mathcal{O}_1 \circ \mathcal{O}_2 \circ \dots \circ \mathcal{O}_n$

Constitutive models can be formalized as **parametrized functions** (stress update):

$$\sigma_{n+1}, \alpha_{n+1} = \mathcal{M}_{\theta}(\varepsilon_{n+1}, \alpha_n)$$

where  $\theta$  are the model **material parameters** and  $\mathcal{M}_{\theta} = \mathcal{O}_1 \circ \mathcal{O}_2 \circ \dots \circ \mathcal{O}_n$

**Formally equivalent with ML**, main **differences** are:

- **internal state** not **observable** from data
- **sequential evaluation of the model** in time:

$$\sigma_1, \alpha_1 = \mathcal{M}_{\theta}(\varepsilon_1, \alpha_0) \rightarrow \sigma_2, \alpha_2 = \mathcal{M}_{\theta}(\varepsilon_2, \alpha_1) \rightarrow \dots \rightarrow \sigma_{n+1}, \alpha_{n+1} = \mathcal{M}_{\theta}(\varepsilon_n, \alpha_n)$$

- some **elementary operations**  $\mathcal{O}_i$  may not have **closed-form** expressions

Constitutive models can be formalized as **parametrized functions** (stress update):

$$\sigma_{n+1}, \alpha_{n+1} = \mathcal{M}_\theta(\varepsilon_{n+1}, \alpha_n)$$

where  $\theta$  are the model **material parameters** and  $\mathcal{M}_\theta = \mathcal{O}_1 \circ \mathcal{O}_2 \circ \dots \circ \mathcal{O}_n$

**Formally equivalent with ML**, main **differences** are:

- **internal state** not **observable** from data
- **sequential evaluation of the model** in time:

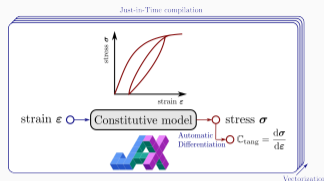
$$\sigma_1, \alpha_1 = \mathcal{M}_\theta(\varepsilon_1, \alpha_0) \rightarrow \sigma_2, \alpha_2 = \mathcal{M}_\theta(\varepsilon_2, \alpha_1) \rightarrow \dots \rightarrow \sigma_{n+1}, \alpha_{n+1} = \mathcal{M}_\theta(\varepsilon_n, \alpha_n)$$

- some **elementary operations**  $\mathcal{O}_i$  may not have **closed-form** expressions

## Opportunities

- **phenomenological** or **NN-based** models can be treated in a **unifying framework**
- they can be **seamlessly hybridized!**
- material parameter **identification** can benefit from **ML gradient-based algorithms**
- **practical implementation** can benefit from **modern ML frameworks**

# JAX for constitutive modeling



**JAX** = accelerated (GPU) array computation and program transformation designed for **HPC** and large-scale **machine learning**

```
def constitutive_update(eps, state, dt):  
    [...]  
    return stress, new_state
```

- **JIT** and **automatic vectorization**

```
batch_constitutive_update = jax.jit(jax.vmap(constitutive_update, in_axes=(0, 0, None)))
```

- **Automatic Differentiation**

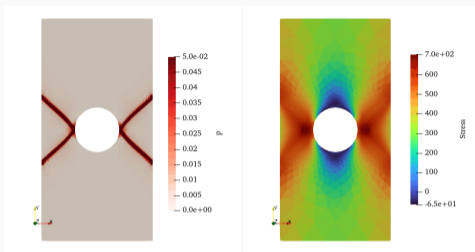
```
constitutive_update_tangent = jax.jacfwd(constitutive_update, argnums=0, has_aux=True)
```

First release end of 2025: <https://github.com/bleyerj/jaxmat>

## Features

- based on JAX and equinox (NN models), optimistix (optimization solvers), diffrax (ODE solvers)
- models as **PyTrees**  $\Rightarrow$  allows **modularity** and **hierarchical composition**:  
 $\theta = \{\theta_{\text{elastic}}, \{\theta_{\text{plastic}}, \theta_{\text{hardening}}\}\}$
- prototype **material models**: elastoplasticity, viscoplasticity, hyperelasticity, finite strain plasticity, viscoelasticity, Generalized Standard Materials, damage, etc.
- **hybrid NN-phenomenological** models
- **differentiable solvers** using **implicit differentiation**
- **gradient-based** identification
- **CPU** or **GPU** acceleration
- **compatibility** with FEniCSx (v.0.10)

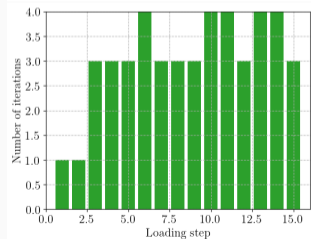
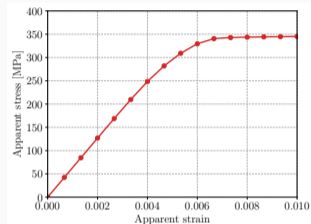
# Small-strain elastoplasticity



```
E, nu = 70e3, 0.3
elastic_model = LinearElasticIsotropic(E, nu)

class VoceHardening(equinox.Module):
    sig0 = 350.0
    sigu = 500.0
    b = 1e3
    def __call__(self, p):
        return self.sig0 + (self.sigu - self.sig0) * (1 - jnp.exp(-self.b *
            p))

material = vonMisesIsotropicHardening(elastic_model,
    yield_stress)
```

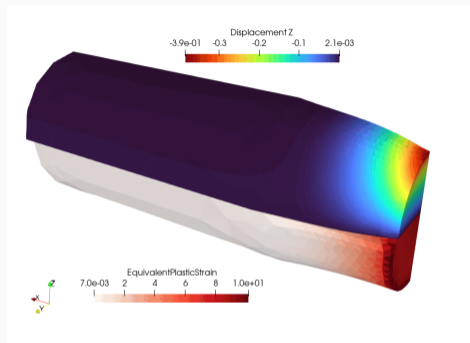


$$\mathbf{F} = \mathbf{F}^e \mathbf{F}^p; \quad \bar{\mathbf{b}}^e = J^{-2/3} \mathbf{F}^e (\mathbf{F}^e)^T$$

$$\boldsymbol{\tau} = \mu \operatorname{dev}(\bar{\mathbf{b}}^e) + \frac{\kappa}{2} (J^2 - 1) \mathbf{I}$$

$$f(\bar{\mathbf{b}}^e) = \mu \|\mathbf{s}\| - \sqrt{\frac{2}{3}} R(p_n + \Delta p) \leq 0$$

$$0 = \operatorname{dev}(\bar{\mathbf{b}}^e - \bar{\mathbf{b}}_{\text{trial}}^e) + \sqrt{\frac{2}{3}} \Delta p \operatorname{tr}(\bar{\mathbf{b}}^e) \frac{\mathbf{s}}{\|\mathbf{s}\|}$$

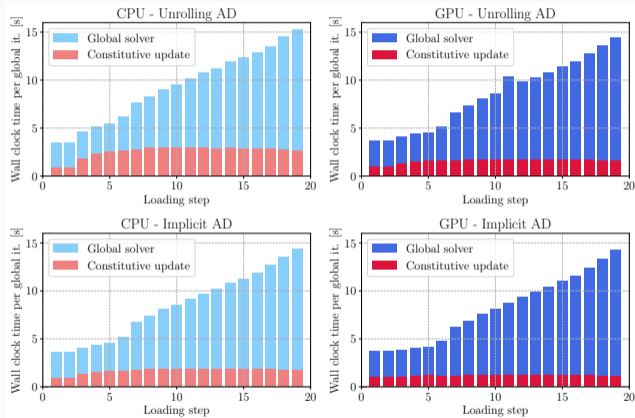


Resolution involves local solving of a Newton system of size 7  
Tangent operator in  $PK1/F$  using **implicit AD**:

$$\mathbf{P} = \boldsymbol{\tau} \mathbf{F}^{-1} \quad \mathbb{C}_{\text{tang}} = \frac{\partial \mathbf{P}}{\partial \mathbf{F}}$$

Constitutive equation: `jax[cpu]` or `jax[gpu]` on NVIDIA RTX A1000

Linear solver: PETSc gmres + gamg



Global linear solves – Constitutive behavior integration

## Identification

$$\theta^* = \arg \min_{\theta} \frac{1}{M} \sum_{k=1}^M \|\sigma^{(k)} - \sigma_{\text{data}}^{(k)}\|^2$$

**gradient-based** optimisation, needs **material parameters sensitivities**

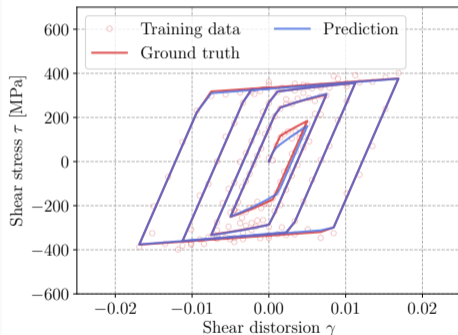
$$\frac{\partial \sigma^{(k)}}{\partial \theta} \rightarrow \frac{\partial F}{\partial \theta}(\varepsilon, \alpha; \theta)$$

easy to obtain with **JAX**

```
@jax.jit
def loss(material, sig_data, loading):
    sig_hat = compute_evolution(material, loading)
    loss = optax.l2_loss(sig_hat, sig_data)
    return loss.mean()
```

differentiable time loop with `jax.lax.scan`

Armstrong-Frederick type hardening plasticity with **noisy data**

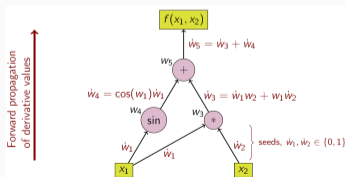




**Thank you for your attention!**

# What is Automatic Differentiation ?

- **Numerical Differentiation:**  $f'(x) = \frac{f(x+h) - f(x)}{h}$  with e.g.  $h = 10^{-6}$   
**truncation/rounding errors,  $O(dim)$  evaluations**
- **Symbolic differentiation:**  $f$  represented as an **expression graph**, generates **another expression graph** of the derivative  
**expression swell, duplicate operations, no-closed form expression**
- **Automatic differentiation:** operates directly on the **computer program**, no symbolic representation (numerical evaluation only), **exact forward and reverse mode (back-propagation in ML)**



[Wikipedia]

```
def f(x):  
    [...]  
    for i in range(n):  
        [...]
```

```
def f(x):  
    if cond:  
        [...]  
    else:  
        [...]
```

```
def f(x):  
    [...]  
    while cond:  
        [...]
```

## Return mapping algorithm

Elastic predictor  $\boldsymbol{\sigma}_{\text{elas}} = \boldsymbol{\sigma}_n + \mathbb{C} : \Delta \boldsymbol{\varepsilon}$

$f_{\text{elast}} = \sigma_{\text{eq}} - R(p_n)$

- if  $f_{\text{elas}} < 0$ :  $\boldsymbol{\sigma}_{n+1} = \boldsymbol{\sigma}_{\text{elas}}$  and  $\Delta p = 0$
- else:  $\boldsymbol{\sigma}_{n+1} = \boldsymbol{\sigma}_{\text{elas}} - 2\mu\Delta\boldsymbol{\varepsilon}^p$  with  $\Delta\boldsymbol{\varepsilon}^p = \Delta p \frac{3}{2\sigma_{\text{eq}}^{\text{elas}}} \mathbf{s}_{\text{elas}}$

$$\text{Solve } r(\Delta p) = \sigma_{\text{eq}}^{\text{elas}} - 3\mu\Delta p - R(p_n + \Delta p) = 0 \quad (1)$$

implement **Solve** using fixed-point algorithm, Newton method, bisection, etc.

**Every step is differentiable** with AD *a priori*, **except** (1).

## Return mapping algorithm

Elastic predictor  $\sigma_{\text{elas}} = \sigma_n + \mathbb{C} : \Delta \epsilon$

$f_{\text{elast}} = \sigma_{\text{eq}} - R(p_n)$

- if  $f_{\text{elas}} < 0$ :  $\sigma_{n+1} = \sigma_{\text{elas}}$  and  $\Delta p = 0$
- else:  $\sigma_{n+1} = \sigma_{\text{elas}} - 2\mu\Delta\epsilon^{\text{P}}$  with  $\Delta\epsilon^{\text{P}} = \Delta p \frac{3}{2\sigma_{\text{eq}}^{\text{elas}}} \mathbf{s}_{\text{elas}}$

$$\text{Solve } r(\Delta p) = \sigma_{\text{eq}}^{\text{elas}} - 3\mu\Delta p - R(p_n + \Delta p) = 0 \quad (1)$$

implement **Solve** using fixed-point algorithm, Newton method, bisection, etc.

**Every step is differentiable** with AD *a priori*, **except** (1).

## Algorithm unrolling

Any algorithm used to solve (1) can be written in JAX using loops, conditionals, etc. We can **differentiate through the algorithm** (*unrolling the algorithm iterations*).

We can leverage instead the **implicit function theorem**

e.g. root finding: Find  $x_\theta$  s.t.  $F(x_\theta; \theta) = 0$

We can leverage instead the **implicit function theorem**

e.g. root finding: Find  $x_\theta$  s.t.  $F(x_\theta; \theta) = 0$

To find  $\partial_\theta x_\theta$ , we differentiate the equation so that:

$$\begin{aligned} [\partial_x F] \partial_\theta x_\theta + \partial_\theta F &= 0 \\ \Rightarrow \partial_\theta x_\theta &= -[\partial_x F]^{-1} \partial_\theta F \end{aligned}$$

need only to solve a **linear system** for the **jacobian matrix**  $[\partial_x F]$

the derivative computation becomes **independent from the algorithm** used to solve the nonlinear system, can use AD to form the jacobian  $[\partial_x F]$

We can leverage instead the **implicit function theorem**

e.g. root finding: Find  $x_\theta$  s.t.  $F(x_\theta; \theta) = 0$

To find  $\partial_\theta x_\theta$ , we differentiate the equation so that:

$$\begin{aligned} [\partial_x F] \partial_\theta x_\theta + \partial_\theta F &= 0 \\ \Rightarrow \partial_\theta x_\theta &= -[\partial_x F]^{-1} \partial_\theta F \end{aligned}$$

need only to solve a **linear system** for the **jacobian matrix**  $[\partial_x F]$

the derivative computation becomes **independent from the algorithm** used to solve the nonlinear system, can use AD to form the jacobian  $[\partial_x F]$

Can be implemented for any function using `jax.lax.custom_root`

Available in `optimistix`:

```
solver = optimistix.Newton(rtol=1e-8, atol=1e-8)
solution = optimistix.root_find(fun, solver, y0, adjoint=optimistix.ImplicitAdjoint())
```

# PANN models for data-driven hyperelasticity

## Isotropic hyperelastic potential

$$\psi(\mathbf{F}) = \Psi(I_1, I_2, I_3)$$

**polyconvexity** guaranteed if  $\Psi$  is **convex**

**Physics-Augmented Neural Networks** [Linden et al., 2023] approximate

$$\psi_{\theta}(\mathbf{F}) = \Psi_{\theta}(I_1, I_2, I_3, -2J)$$

where  $\Psi_{\theta}(\mathcal{I})$  is an **Input Convex Neural Network** (ICNN) [Amos et al., 2017] (positive weights, softplus activation)

**Important:** enforce stress-free initial configuration

$$\Psi_{\theta}^{\text{PANN}}(\mathbf{C}) = \Psi_{\theta}(\mathcal{I}) - \left. \frac{\partial \Psi_{\theta}}{\partial \mathcal{I}} \right|_{\mathbf{C}=\mathbf{I}} 2(J - 1)$$

```
class PANN(ICNN):
    def nn_energy(self, lambC):
        I3 = jnp.prod(lambC)
        I1 = jnp.sum(lambC)
        I2 = jnp.sum(I3 / lambC)
        return super().__call__(jnp.asarray([I1, I2,
        I3, -2*jnp.sqrt(I3)]))
```

```
def pann_energy(self, lambC):
    J = jnp.sqrt(jnp.prod(lambC))
    dpsidC =
    jax.jacfwd(self.nn_energy)(jnp.ones((3,))) [0]
    return self.nn_energy(lambC) - 2 * dpsidC *
    (J - 1)
```

# PANN model on rubber elasticity data [Treloar, 1944]

PANN with 4 inputs, 10 neurons in 1 hidden layer

